

Structures and Bonding in  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$ Hiroshi Mizoguchi,<sup>†</sup> Danielle Gray,<sup>†</sup> Fu Qiang Huang,<sup>†,‡</sup> and James A. Ibers<sup>\*,†</sup>

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The compounds  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$ , members of the  $AAn_2Q_6$  actinide family ( $A$  = alkali metal or Tl;  $An$  = Th or U;  $Q$  = S, Se, or Te), have been synthesized from  $US_2$ ,  $K_2S$ , and S at 1273 K and U,  $K_2Se$ , and Se at 1173 K, respectively.  $KU_2Se_6$  shows Curie–Weiss behavior above 30 K and no magnetic ordering down to 5 K. The value of  $\mu_{\text{eff}}$  is 2.95(1)  $\mu_B/U$ . Its electronic spectrum shows the peaks characteristic of 5f–5f transitions. It is a semiconductor with an activation energy of 0.27 eV for electrical conduction. Both  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$  crystallize in space group  $Immm$  of the orthorhombic system and are of the  $KTh_2Se_6$  structure type. Both contain infinite one-dimensional linear Q–Q chains characteristic of the  $AAn_2Q_6$  family. Typical of the known  $AAn_2Q_6$  compounds, in  $KU_2Se_6$ , there are two alternating Se–Se distances of 2.703(2) and 2.855(2) Å, both much longer than an Se–Se single bond. In contrast, in  $K_{0.91}U_{1.79}S_6$ , the first sulfide of this family to be characterized structurally, there are alternating normal  $S_2^{2-}$  pairs 2.097(5) Å in length. In  $K_{0.91}U_{1.79}S_6$ , the formal oxidation state of U is 4+.

## Introduction

The actinide compounds of formula  $AAn_2Q_6$  ( $A$  = alkali metal or Tl;  $An$  = Th or U;  $Q$  = S, Se, or Te) present interesting problems in chemical bonding. When  $Q$  = Te, the known compounds  $CsTh_2Te_6$ ,<sup>1</sup>  $Tl_{1.12}U_2Te_6$ ,<sup>2</sup> and  $KTh_2Te_6$ ,<sup>3</sup> crystallize in space group  $Cmcm$  of the orthorhombic system. These compounds are isostructural; the layered structure features not only isolated  $Te^{2-}$  species but linear, infinite, one-dimensional Te chains with Te–Te distances approximately 0.35 Å longer than that of a typical Te–Te single bond. When  $Q$  = Se, the known compounds  $KTh_2Se_6$ ,<sup>4</sup>  $RbTh_2Se_6$ ,<sup>4</sup> and  $CsU_2Se_6$ ,<sup>5</sup> crystallize in space group  $Immm$  of the orthorhombic system. These compounds, which are isostructural, have a layered structure that differs in its stacking pattern from that of the  $Cmcm$  structure. The  $Immm$  structure contains discrete  $Se^{2-}$  species and infinite one-

dimensional Se chains with Se–Se distances about 0.35 Å longer than a typical Se–Se single bond.

Efforts have been made to understand the bonding in these compounds when  $Q$  = Se. Thus, it has been postulated that in the  $ATh_2Se_6$  ( $A$  = K, Rb) compounds the Se atoms in the  $[Th_2Se_6]$  layers accept the extra electrons from the A atoms, thereby breaking one out of four  $Se_2^{2-}$  bonds.<sup>4</sup> This picture is supported by a variety of physical measurements on these compounds.<sup>4</sup> This formulation leaves Th with the expected formal oxidation state of 4+. On the basis of XPS measurements, a formal oxidation state of 4+ was assigned to U in  $KU_2Se_6$ .<sup>5</sup> In an effort to clarify further the bonding in the  $AAn_2Q_6$  compounds, we report here the syntheses, structure, and some physical properties of  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$ .

## Experimental Section

**Synthesis of  $K_{0.91}U_{1.79}S_6$ .** The compound was prepared by the reactive flux method.<sup>6</sup>  $K_2S$  was prepared by the stoichiometric reaction of K (Cerac, 98%) and S (Mallinckrodt, 99.6%) in liquid  $NH_3$  at 194 K.  $\beta$ - $US_2$  was prepared by the stoichiometric reaction of U turnings (depleted, ORNL) with S in a carbon-coated fused-silica ampule. The heating profile was 673 K for 20 h, then 873 K for 20 h, followed by slow cooling to 298 K.

A reaction mixture of  $K_2S$  (0.5 mmol),  $\beta$ - $US_2$  (0.5 mmol), and S (0.5 mmol) was loaded into a fused-silica ampule in an Ar-filled

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glovebox. The ampule was evacuated to  $10^{-4}$  Torr and sealed. The sample was heated to 1273 K in 23 h, kept at 1273 K for 48 h, cooled at 10 K/h to 773 K, and then cooled at 120 K/h to 298 K. Black plates of the desired product  $K_{0.91}U_{1.79}S_6$  were obtained in about 5% yield, based on U. These plates were manually extracted from the melt.

**Synthesis of  $KU_2Se_6$ .** An earlier preparation<sup>5</sup> of  $KU_2Se_6$  involved the reaction of  $K_2Se$ , U, and Se at 773 K. Here the reaction mixture consisted of  $K_2Se$  (0.17 mmol) (prepared by the stoichiometric reaction of K and Se (Cerac, 99.99%) in liquid  $NH_3$  at 194 K), U (0.5 mmol), and Se (1.4 mmol). The sample was heated to 1173 K in 15 h, held at 1173 K for 96 h, cooled at 5.75 K/h to 603 K, and then cooled at 3 K/h to 298 K. The yield of golden-black needles of  $KU_2Se_6$  was approximately 25%, based on U.

Selected single crystals of each of the above compounds were examined with an EDX-equipped Hitachi S-3500 SEM. The presence of the stated elements in approximately 1:2:6 ratios was confirmed.

**Syntheses of  $USe_3$  and  $US_3$ .**  $USe_3$  was prepared from the elements.<sup>7</sup> Its X-ray powder pattern was consistent with that calculated from single-crystal data.<sup>7</sup>  $US_3$  was also prepared from the elements. Its powder pattern agreed with one given previously.<sup>8</sup>

**Structure Determinations.** Single-crystal X-ray diffraction data were collected by the use of graphite-monochromated Mo  $K\alpha$  radiation ( $\lambda = 0.71073$  Å) at 153 K on a Bruker SMART-1000 CCD diffractometer.<sup>9</sup> The crystal-to-detector distance was 5.023 cm. Data were collected by scans of  $0.3^\circ$  in  $\omega$  in four groups of 606 frames at  $\varphi$  settings of  $0^\circ$ ,  $90^\circ$ ,  $180^\circ$ , and  $270^\circ$  at exposure times of 15 s/frame. The collection of the intensity data was carried out with the program SMART.<sup>9</sup> Crystal decay was monitored by re-collecting 50 initial frames at the end of each data collection. Cell refinement and data reduction were carried out with the use of the program SAINT+,<sup>9</sup> and face-indexed absorption corrections were performed numerically with the use of the program XPREP.<sup>10</sup> Because the crystals were so thin, a Leitz microscope with a calibrated traveling micrometer eyepiece was employed to measure accurately the dimensions of the crystal faces for the face-indexed absorption corrections. Then, the program SADABS<sup>10</sup> was used to make incident beam and decay corrections.

Each crystal structure was solved with the direct methods program SHELXS and refined with the full-matrix least-squares program SHELXL.<sup>10</sup> The program STRUCTURE TIDY<sup>11</sup> was used to standardize the positional parameters. After refinement of the standard anisotropic model including extinction, the occupancy of the U site was fixed but the occupancies of the alkali metal A and the chalcogen Q sites were varied in an ensuing refinement. The resultant formulas were  $K_{1.01(2)}U_{2}S_{6.69(6)}$  and  $K_{0.98(2)}U_{2}Se_{6.06(2)}$ . The Q = Se formula does not differ significantly from the stoichiometric formula  $KU_2Se_6$ , and in the final refinement, the site occupancies were fixed at this composition. On the other hand, the Q = S compound is clearly nonstoichiometric, the above formula being equivalent to  $K_{0.91(2)}U_{1.79(2)}S_6$ , indicative of K and U vacancies. Accordingly, in a final refinement, the occupancies of the two S sites were fixed but the occupancies of the U and K sites were varied. This refinement converged to give the same formula, namely

**Table 1.** Crystal Data and Structure Refinements for  $KU_2Q_6^a$

	$K_{0.91}U_{1.79}S_6$	$KU_2Se_6$
$a$ , Å	4.0321(4)	4.0586(6)
$b$ , Å	5.3929(5)	5.5578(8)
$c$ , Å	20.556(2)	21.683(3)
$V$ , Å <sup>3</sup>	446.99(7)	489.1(1)
$\rho_c$ , g/cm <sup>3</sup>	4.856	6.715
$\mu$ , cm <sup>-1</sup>	341.3	557.4
$R(F)^b$	0.0248	0.0204
$R_w(F_o^2)^c$	0.0582	0.0564

<sup>a</sup> For both structures,  $Z = 2$ , space group =  $Immm$ ,  $T = 153$  K, and  $\lambda = 0.71073$  Å. <sup>b</sup>  $R(F) = \sum ||F_o| - |F_c||/|F_o|$  for  $F_o^2 > 2\sigma(F_o^2)$ . <sup>c</sup>  $R_w(F_o^2) = \{\sum [w(F_o^2 - F_c^2)^2]/\sum wF_o^4\}^{1/2}$  for all data.  $w^{-1} = \sigma^2(F_o^2) + (q_1P)^2 + q_2P$ , where  $P = (\max(F_o^2, 0) + 2F_c^2)/3$ .  $q_1 = 0.0156$ ,  $q_2 = 22.25$  for  $KU_2S_6$ ; 0.0280, 0 for  $KU_2Se_6$ .

$K_{0.91(2)}U_{1.79(2)}S_6$ . Although the determination of composition from the refinement of X-ray occupancies is not without its pitfalls,<sup>12</sup> there are two features that lead us to believe that the present calculations are reliable. The first is that the structure is a simple one that involves only one U, one K, and two Q positions in the asymmetric unit. Whether one varied the occupancies of K and U while keeping the occupancies of Q1 and Q2 fixed or whether one varied the occupancies of K, Q1, and Q2 while keeping the occupancy of U fixed, the same results were obtained. Thus, the refinement is stable and convergence to a false minimum is highly unlikely. The second is that in each refinement the ratio Q1:Q2, which is 1:2 for the stoichiometric structure, remained at 1:2 within statistical error. Additional crystallographic details and selected metrical data are given in Table 1 and in the Supporting Information.

The unit cell parameters for  $KU_2Se_6$  of 4.0586(6), 5.5578(8), and 21.683(3) Å ( $a$ ,  $b$ ,  $c$ ) at 153 K in Table 1 may be compared with those reported earlier<sup>5</sup> of 4.058(3), 5.556(4), and 21.710(17) Å at 298 K. (No structure determination was carried out in the earlier investigation.)

**Magnetic Susceptibility Measurement.** The magnetic susceptibility data were collected for a 22 mg sample of  $KU_2Se_6$  in the temperature range 5–390 K with a measuring field of 500 G, after cooling in either a zero field or at 10 kG. The crushed sample was contained in a gelatine capsule. A Quantum Design MPMS-5 SQUID magnetometer was used. All measurements were corrected for core diamagnetism.<sup>13</sup>

**Diffuse Reflectance Measurements.** Diffuse reflectance data for  $KU_2Se_6$ ,  $USe_3$ , and  $US_3$  were collected over the spectral range 500–2600 nm with a Varian Cary 5000 double-beam spectrometer. MgO was used as a reference. The Kubelka–Munk function was used to transform the data into absorbance.

**Electrical Conductivity Measurements.** Standard four-probe electrical conductivity measurements were made over the temperature range 200–300 K on a plate-shaped single crystal of  $KU_2Se_6$ . Electrical contacts were applied with the use of silver paste. A Quantum Design PPMS instrument was used.

## Results and Discussion

**Syntheses.** The synthesis of  $K_{0.91}U_{1.79}S_6$  proceeded from the reaction of  $K_2S$ ,  $US_2$ , and S at 1273 K. It is an inefficient synthesis, affording only about a 5% yield of black plates of the desired product. Attempts to increase the yield and the size of the crystals obtained were unsuccessful. On the

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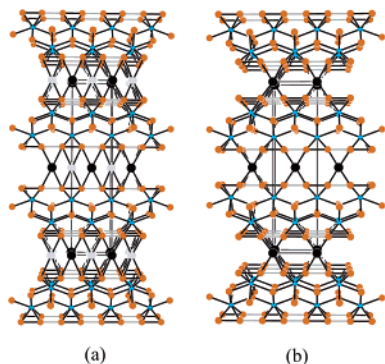
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**Figure 1.** Structures of (a)  $CsU_2Se_6$  and (b)  $KU_2Se_6$  as viewed down the  $a$  axis. Color code: orange = Se; blue = U; black = A; gray = disordered position of A; black bond = shorter distance in the Se–Se chain; gray bond = longer distance.

other hand, the synthesis of  $KU_2Se_6$  from  $K_2Se$ , U, and Se at 1173 K afforded high-quality golden-black needles in about 25% yield. The compounds are stable in air for a few hours.

**Structures.** The two structure types comprising the known  $AAn_2Q_6$  structures are those of  $CsTh_2Te_6$  (space group  $Cmcm$ )<sup>1</sup> and the closely related  $KTh_2Se_6$  (space group  $Immm$ ).<sup>4</sup> Both are layered structures. The difference between the  $Cmcm$  and  $Immm$  structure types lies in the stacking of the  $UQ_3$  layers and the positioning of the A atom between the layers. In the  $Cmcm$  structure, each succeeding  $UQ_3$  layer is shifted by  $1/2a$ . In the  $Immm$  structure, each succeeding  $UQ_3$  layer is shifted by  $1/2a$  and  $1/2b$ .

$CsU_2Se_6$ ,<sup>5</sup>  $K_{0.91}U_{1.79}S_6$ , and  $KU_2Se_6$  crystallize with the  $KTh_2Se_6$  structure type in space group  $Immm$ . However, the Cs atom is disordered over two positions in the structure of  $CsU_2Se_6$ ,<sup>5</sup> whereas the K atoms are ordered in the structures of  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$ . This is a steric effect. The structures of  $CsU_2Se_6$ <sup>5</sup> and  $KU_2Se_6$  are shown in Figure 1. In addition to isolated Q atoms in the  $UQ_3$  layers, linear chains of Q atoms with alternating Q–Q distances are found in these structures. The details of the linear Q chains characteristic of the  $AAn_2Q_6$  family and related compounds bear on the question of chemical bonding, and these are summarized in Table 2. In the structure of  $CsU_2Se_6$ ,<sup>5</sup> disorder occurs because the Cs atom can either occupy the space between four short Se–Se pairs or the space between four long Se–Se pairs. In each of the  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$  structures, the K atom only occupies the position between the four long Q–Q pairs because it will not physically fit in the space between the short Q–Q pairs.

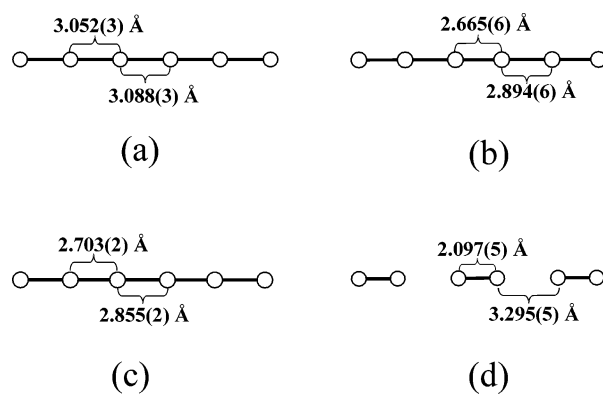
In  $KU_2Se_6$ , the two alternating Se–Se distances of 2.703(2) and 2.855(2) Å are both considerably longer than a Se–Se single bond distance, for example, that of 2.363(1) Å in  $USe_3$ ;<sup>7</sup> in  $K_{0.91}U_{1.79}S_6$ , the S–S distances are 2.097(5) Å (a single bond distance) and 3.295(5) Å (a nonbonding distance). The Q–Q chains in these and related compounds are compared in Figure 2.

Selected distances for  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$  are presented in Table 3. These are normal. The U–Q distances, which range from 2.761(2) to 2.846(2) Å for  $K_{0.91}U_{1.79}S_6$ ,

**Table 2.** Q–Q Distances in the Infinite Linear Q–Q Chains of Selected Compounds

compound	$d_{short}$ (Å)	$d_{long}$ (Å)	$d_{short}/d_{long}$	ref
$K_{0.91}U_{1.79}S_6^a$	2.097(5)	3.295(5)	0.636	this work
$KU_2Se_6^a$	2.703(2)	2.855(2)	0.947	this work
$CsU_2Se_6^a$	2.665(6)	2.894(6) <sup>d</sup>	0.921	5
$Tl_{1.12}U_2Te_6^{b,c}$	3.037(3)	3.054(3)	0.994	2
$KTh_2Se_6^a$	2.727(2)	2.907(2)	0.938	4
$RbTh_2Se_6^a$	2.728(3)	2.906(3)	0.939	4
$KTh_2Te_6^b$	3.057(3)	3.085(3)	0.991	3
$CsTh_2Te_6^b$	3.052(3)	3.088(3)	0.988	1
$US_3$	2.086(4)	3.290(4)	0.634	17
$USe_3$	2.363(1)	3.294(5) <sup>d</sup>	0.717	7
$\alpha$ - $UTe_3$	2.751(1)	3.350(1)	0.821	29
$ZrS_3$	2.090(5)	3.035(5)	0.689	30
$ZrSe_3$	2.34	3.07	0.76	31
$ZrTe_3$	2.794(1)	3.105(1)	0.900	32

<sup>a</sup> Space group  $Immm$ . <sup>b</sup> Space group  $Cmcm$ . <sup>c</sup> A more extensive tabulation of Te–Te linear chain geometries is given in ref 33. <sup>d</sup> Distance not given in the reference but calculated from data therein.



**Figure 2.** Q–Q chains in (a)  $CsTh_2Te_6$ , (b)  $CsU_2Se_6$ , (c)  $KU_2Se_6$ , and (d)  $K_{0.91}U_{1.79}S_6$ .

**Table 3.** Selected Distances (Å) for  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$

	$K_{0.91}U_{1.79}S_6$	$KU_2Se_6$
K–Q(2) × 8	3.339(2)	3.3574(6)
U–Q(1) × 2	2.761(2)	2.8884(7)
U–Q(1) × 2	2.8082(8)	2.9648(5)
U–Q(2) × 4	2.846(2)	2.9241(6)
Q(2)–Q(2)	2.097(5)	2.703(2)

may be compared with those of 2.76–2.99 Å for  $UFeS_3$ ;<sup>14</sup> the U–Se distances, which range from 2.8884(7) to 2.9648(5) Å for  $KU_2Se_6$ , may be compared with those of 2.850(2), 2.935(2), and 3.126(3) Å for  $UMnSe_3$ .<sup>15</sup> The K–Q distance of 3.339(2) Å for  $K_{0.91}U_{1.79}S_6$  is in the range of 3.127(2)–3.382(2) Å found for  $KCuZrS_3$ ;<sup>16</sup> the K–Se distance of 3.3574(6) Å may be compared with those of 3.255(2)–3.530(2) Å for  $KCuZrSe_3$ .<sup>16</sup>

The structure of  $KU_2Se_6$  may be thought of as resulting from the insertion of K atoms into the  $USe_3$  structure.<sup>7</sup> Such insertion weakens the Se–Se bonds, either by formal oxidation of the K atom or by reduction of the U atom. The stoichiometric formula  $KU_2Se_6$  established from the refinement of the X-ray data does not distinguish these two

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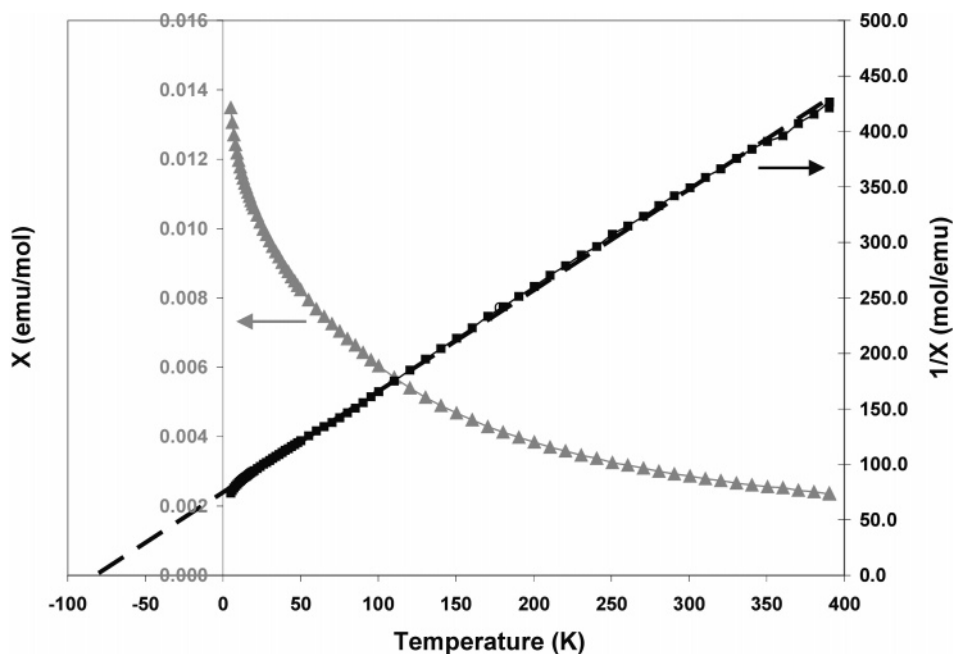


Figure 3. Magnetic susceptibility,  $\chi$ , and  $\chi^{-1}$  vs  $T$  for  $\text{KU}_2\text{Se}_6$ .

possibilities. The formal oxidation state of U in  $\text{KU}_2\text{Se}_6$  was determined to be 4+ from XPS measurements,<sup>5</sup> although it is not clear if the method is sufficiently sensitive to eliminate other possible oxidation states in view of the breadth of the binding energy signal and the possibility of surface contamination. As opposed to  $\text{KU}_2\text{Se}_6$ , in  $\text{K}_{0.91}\text{U}_{1.79}\text{S}_6$ , the S–S single bonds are not weakened. Given the composition  $\text{K}_{0.91(2)}\text{U}_{1.79(2)}\text{S}_6$ , as established from the refinement of the X-ray data, the stoichiometric formulation  $\text{KU}_2(\text{S})_2(\text{S}_2)_2$  as indicated from the structure, and charge balance, the formal oxidation state of U in this compound is 3.96(5)+, hence 4+. From single-crystal diffraction studies, both  $\text{US}_3$ <sup>17</sup> and  $\text{USE}_3$ <sup>7</sup> are of the  $\text{ZrSe}_3$  structure type. If we think of  $\text{K}_{0.91(2)}\text{U}_{1.79(2)}\text{S}_6$  as resulting from the insertion of K into the  $\text{US}_3$  structure with no change in the  $\text{U}^{4+}$  oxidation state, then we can envision the series  $\text{K}_x\text{U}_{2-x/4}\text{S}_6$ ,  $1 \geq x \geq 0$ , as we progress from  $\text{KU}_{1.75}\text{S}_6$  to  $\text{U}_2\text{S}_6$  ( $\text{US}_3$ ). In this series, the number of metal vacancies is  $1 - 0.75x$ , so clearly at some point as  $x$  decreases the  $\text{KU}_{1.75}\text{S}_6$  structure will collapse. In the present instance of  $x = 0.91$ , the structure is stable, but undoubtedly other values of  $x$  are possible and may depend on reaction conditions.

**Magnetic Susceptibility.** The magnetic susceptibility,  $\chi$ , and its inverse,  $\chi^{-1}$ , as a function of temperature for  $\text{KU}_2\text{Se}_6$  are shown in Figure 3. The field-cooled and zero-field-cooled data are essentially superimposable. Above 30 K, the Curie–Weiss law,  $\chi = C/(T - \theta)$ , is obeyed. The Curie constant and the Weiss temperature obtained from parameter fitting are 2.17(1)  $\text{emu}\cdot\text{K}/\text{mol}$  and  $-81.2(6)$  K, respectively. The effective magnetic moment,  $\mu_{\text{eff}}$ , calculated from the Curie constant is 2.95(1)  $\mu_{\text{B}}/\text{U}$ . The values of  $\mu_{\text{eff}}$  for the free ions  $\text{U}^{3+}$  ( $6d^05f^3$ ),  $\text{U}^{4+}$  ( $6d^05f^2$ ), and  $\text{U}^{5+}$  ( $6d^05f^1$ ) calculated using L–S coupling are 3.62, 3.58, and 2.54  $\mu_{\text{B}}/\text{U}$ ,

U, respectively. The observed value of 2.95(1)  $\mu_{\text{B}}/\text{U}$  in  $\text{KU}_2\text{Se}_6$  is consistent with values previously observed in solid-state chalcogenides in which U has a formal oxidation state of 4+.<sup>18–20</sup> For example, a value of 2.50  $\mu_{\text{B}}/\text{U}$  has been reported for the compound  $\text{Cu}_2\text{U}_3\text{S}_7$  in which the absence of S–S bonds and the propensity of Cu to be 1+ in chalcogenides implies  $\text{U}^{4+}$ .<sup>18</sup> Comparable comparisons are lacking for solid-state chalcogenides in which U has a formal oxidation state of 3+.

Magnetic ordering was not observed in  $\text{KU}_2\text{Se}_6$  until 5 K. The negative Weiss constant implies antiferromagnetic interactions among U ions. The related compound  $\text{USE}_3$  ( $\text{USe}(\text{Se}_2)$ ) exhibits antiferromagnetic-like behavior at 40 K.<sup>21</sup> The  $\text{U}\cdots\text{U}$  distances in  $\text{KU}_2\text{Se}_6$  are 4.059(6) (*ab* plane) and 7.753(1) Å (interplane), whereas they are 4.06(1) (*ab* plane) and 7.02(1) Å (interplane) in  $\text{USE}_3$ .<sup>7</sup> The experimental observation that magnetic ordering is less facile as the interplane  $\text{U}\cdots\text{U}$  distance increases is consistent with the Hill plot.<sup>22</sup>

**Diffuse Reflectance Spectra.** Figure 4 shows the optical spectra for  $\text{KU}_2\text{Se}_6$ ,  $\text{USE}_3$ , and  $\text{US}_3$ . The characteristic peaks from f–f transitions observed for  $\text{USE}_3$  and  $\text{US}_3$  are also observed for  $\text{RbU}_2\text{SbS}_8$  and  $\text{KU}_2\text{SbSe}_8$ <sup>23</sup> at nearly the same positions. All four compounds formally contain  $\text{U}^{4+}$ .<sup>23</sup> Because of the interference of these characteristic peaks, the band gaps of these materials can only be estimated to be in the region 0.4–0.8 eV. The absorptions originating from 5f–

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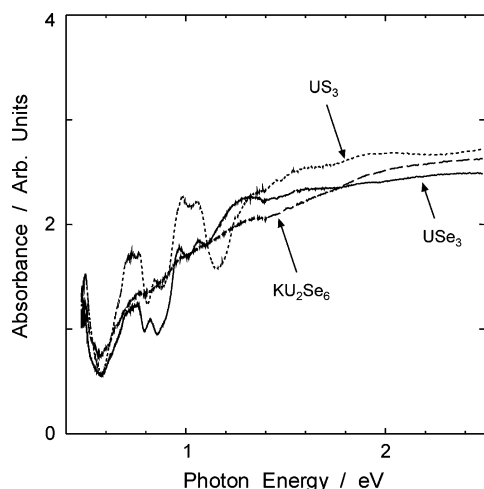
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**Figure 4.** Photon energy vs absorbance for  $KU_2Se_6$ ,  $US_3$ , and  $USe_3$ .

5f transitions are much weaker in  $KU_2Se_6$  than in the other compounds.

**Electrical Conductivity.** The value of the electrical conductivity along the  $a$  axis of a  $KU_2Se_6$  single crystal is  $1 \times 10^{-3}$  S/cm at 300 K. The material is a semiconductor with conductivity decreasing with decreasing temperature in a simple Arrhenius-type thermally activated behavior. The derived activation energy of 0.27 eV for electrical conductivity is consistent with a band gap larger than 0.54 eV and hence consistent with the diffuse reflectance spectrum.

**Electronic Structure.** Why do the Q $\cdots$ Q chains in  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$  differ so markedly (Figure 2)? Ideally, band structure calculations for these compounds would be helpful in answering this question. However, such calculations for U compounds are complicated by a number of factors, including large spin-orbit interactions and electron repulsion for contracted U 5f orbitals, which is Mott Hubbard splitting. Therefore, we offer a qualitative argument. The binding energies for the single-bonded species are 226 kJ/mol<sup>24</sup> for  $S_2^{2-}$  and 209 kJ/mol<sup>24</sup> for  $Se_2^{2-}$ . Although no value for  $Te_2^{2-}$  appears to be known, it should be less than that for  $Se_2^{2-}$ . From the molecular orbital diagram for  $Q_2^{2-}$ ,<sup>25</sup> it is known that near the HOMO-LUMO gap the states are Q np Q np  $\sigma$ ,  $\pi$ ,  $\pi^*$ , and  $\sigma^*$ . The  $\pi$  and  $\pi^*$  states are doubly degenerate. These four states are modified by Q ns Q ns  $\sigma$  bonding. At the HOMO, the  $\pi^*$  states are fully occupied by electrons, indicating the existence of a Q-Q single bond. The LUMO is Q np Q np  $\sigma^*$ . Because of their binding energies, the position of the LUMO of  $S_2^{2-}$  is higher relative to that of  $Se_2^{2-}$ .

We consider solids constructed from such  $Q_2^{2-}$  species. The compound  $USe_3$  contains two different Se $\cdots$ Se distances, 2.361(5) and 3.294(5) Å. The average Se-Se bond length in  $\alpha$ - $Se_8$  at 299 K is 2.336(7) Å,<sup>26</sup> a typical nonbonding Se-

$\cdots$ Se distance is around 3.3 Å, for example, 3.269(2) Å in  $Gd_2Se_3$ .<sup>27</sup> Therefore, for  $USe_3$ , the valence description  $U^{4+}(Se^{2-})(Se_2^{2-})$  is plausible. The electronic structure around the Fermi energy is dominated by the interaction among the  $U^{4+}$  6d5f, Se 4p  $\pi/\pi^*$  (occupied states), and Se 4p  $\sigma^*$  (unoccupied state) for  $Se_2^{2-}$ . This black compound is semiconductive.<sup>28</sup>

Next, we consider  $KU_2Se_6$  where the Se-Se bonds are longer than in  $USe_3$  and discrete Se-Se single-bonded pairs are no longer present (Table 2). The  $Se_2^{2-}$ -species has been reduced, presumably by an electron from the K atom that is inserted into the Se 4p  $\sigma^*$  orbital. From charge balance, this compound is represented as  $KU^{4+}_2(Se^{2-})_2(Se^{1.25-})_4$  with one-fourth of the Se 4p  $\sigma^*$  band occupied by the inserted electron. The  $Se^{1-}$  chain in  $USe_3$  and the  $Se^{1.25-}$  chain in  $KU_2Se_6$  can be derived by adding extra electrons to a linear chain of equally spaced neutral Se atoms. The linear Se chain is distorted by the intrinsic charge density waves and the cation ( $K^+$ ,  $U^{4+}$ ) distribution. Thus, we would expect semiconducting behavior in these compounds, consistent with experiment.

Finally, we consider  $K_{0.91}U_{1.79}S_6$ . The electron responsible for the reduction of the Se-Se single bond in the  $Se_2^{2-}$ -species in  $KU_2Se_6$  cannot reduce the stronger S-S single bond in the  $S_2^{2-}$ -species because of the relatively higher position of the S 3p  $\sigma^*$  orbital compared to the Se 4p  $\sigma^*$  orbital. Electron insertion into the  $S_2^{2-}$  dimer chain leading to a delocalized S chain with S-S distances longer than a normal single bond does not lower the total energy of the system as effectively as the formation of cationic vacancies. The structural difference between the sulfide and the selenide arises because the S valence states are more localized (anionic) than are the Se and Te states. Therefore, a compensation takes place via the formation of cationic vacancies that results in the observed formulation  $K_{0.91}U_{1.79}S_6$ . Apparently, the formation of cationic vacancies is also more facile than reduction of the  $U^{4+}$  ion.

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**Supporting Information Available:** Crystallographic data in CIF format for  $K_{0.91}U_{1.79}S_6$  and  $KU_2Se_6$ . This material is available free of charge via the Internet at <http://pubs.acs.org>.

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